

High-Resolution Microwave Phonon Spectroscopy of Dispersion-Shifted Fiber

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Abstract—First measurements of acoustic modes in dispersion-shifted fiber are presented. High-resolution microwave phonon spectroscopy, performed using an improved method for resolving resonant spectral features, reveals the presence of the L_{04} mode. The use of external cavity tunable lasers for pump and probe enables kilohertz range spectral resolution and high sensitivity. In addition, this method enables down conversion of the phonon spectrum to lower frequencies through the use of beat-brillouin effect.

Index Terms—Acoustic mode, Brillouin amplification, dispersion-shifted fiber, phonon spectroscopy.

I. INTRODUCTION

THE LOW THRESHOLD and narrow bandwidth of Brillouin scattering may be used in applications such as Brillouin amplification for coherent systems and tunable optical filters for wavelength-division-multiplexing (WDM) systems [1], [2]. Experimental methods to resolve the Brillouin spectra in bulk material have been reported both in time and frequency domains [3]–[6]. In optical fibers the spectra can be analyzed in terms of normal acoustic modes [7], [8]. It has been observed that Brillouin spectrum in optical fibers has fine structure resulting from the discrete nature of longitudinal acoustic modes [8], [9]. In the optical heterodyne method, acoustic modes in optical fiber were resolved by tuning the probe frequency of a distributed feedback (DFB) laser [9] with the resolution being limited to the linewidth of the DFB laser. Nikles *et al.*, have proposed a technique to resolve Brillouin spectrum by using a swept microwave side band generated by an electrooptic modulator acting on a single Nd:YAG laser [10]. In this letter, we propose and demonstrate an improvement to this technique that provides not only fine resolution via the use of microwave side band as a probe signal, but also a high sensitivity through the amplification of weak acoustic modes. In addition, the use of external cavity tunable lasers for pump and probe enables down conversion of the phonon spectrum to lower frequencies through the use of the beat-brillouin effect. We also report, for the first time, the acoustic mode spectrum of the dispersion-shifted fiber (DSF) and the first experimental observation of the third higher order acoustic mode of dispersion-shifted optical fiber.

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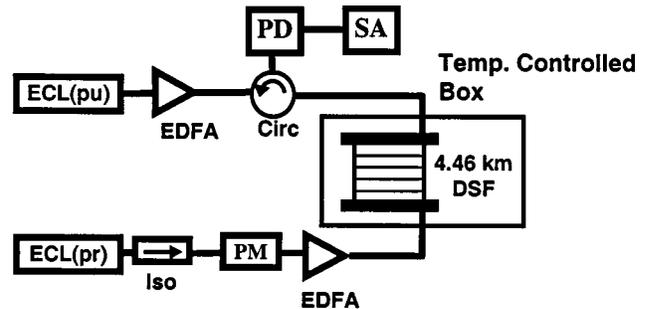


Fig. 1. Experimental setup for the microwave sideband phonon spectroscopy (MSPS). EDFA: Erbium-doped fiber amplifier. PM: Phase modulator. Circ.: Circulator. Iso.: Isolator. PD: Photodetector.

II. PRINCIPLE OF OPERATION

Stimulated Brillouin amplification is a well-known nonlinear effect in optical fibers [1]. A pump wave induces a moving grating by means of the electrostrictive effect in optical fiber. The grating backscatters the pump wave into the Stokes wave, down shifted in frequency from the pump by approximately 11 GHz. When a counterpropagating modulated probe wave is introduced, modulation sidebands that coincide with the bandwidth of the Stokes wave will be amplified [11]. To describe the principle, we refer to the experimental setup in Fig. 1. A temperature-stabilized 4.46-km DSF (Corning SMF/DS) was used as the gain medium. Two External Cavity Lasers (ECL) were used as the pump (ECL-pu) and the probe beam (ECL-pr). The ECLs exhibit linewidth below 150 kHz, as confirmed by the standard heterodyne linewidth measurement technique. The use of an ECL offers us two orders of magnitude better frequency resolution than a DFB laser.

In the optical domain, the backscattered signal contains the following three components: 1) the Fresnel and Rayleigh backscattered pump signal at frequency (ν_{pu}); 2) the Brillouin signal at $\nu_{pu} - \nu_B$, where ν_B is the Stokes shift; and 3) the probe signal at ν_{pr} . After photodetection, the microwave spectrum contains the Brillouin signal centered at ν_B , a beat tone at $(\nu_{pu} - \nu_{pr})$, and a beat-Brillouin tone at $\nu_{pr} - (\nu_{pu} - \nu_B)$. When the phase modulation frequency equals $\nu_{pr} - (\nu_{pu} - \nu_B)$, phase to amplitude conversion occurs resulting in a peak in the measured microwave spectrum. Similarly, an amplitude modulation sideband falling under the Brillouin gain curve results in amplification of the modulation signal. By choosing phase or amplitude modulation, real or imaginary part of the Brillouin spectrum can be mapped.

Fig. 2 shows the microwave spectrum for the beat-Brillouin signal centered at $\nu_{pr} - (\nu_{pu} - \nu_B)$, with and without phase

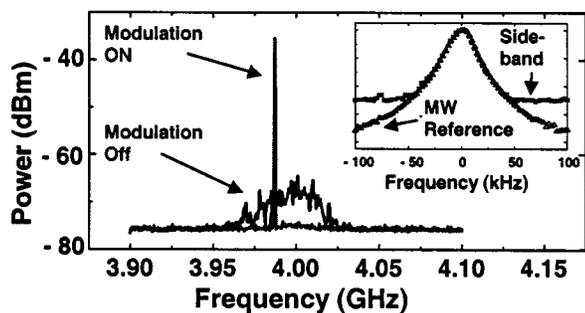


Fig. 2. Brillouin amplification of the modulation sideband. Inset shows the reference signal and the sideband.

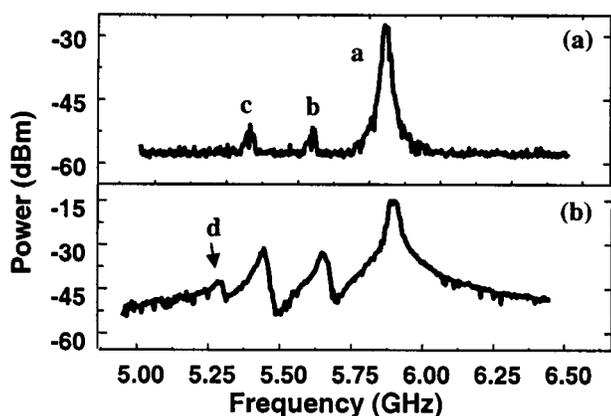


Fig. 3. Measured phonon spectra of DSF. (a) Probe power: 2 dBm. Pump power: 13 dBm. (b) Probe power: 9 dBm. Pump power: 13 dBm.

modulation. The gain spectrum of the DSF was measured to be centered at $\nu_B = 10.5$ GHz with a FWHM, $\Delta\nu_B = 25$ MHz. When the phase modulation is on and its frequency lies within the beat-Brillouin spectrum, a strong narrow linewidth signal, at the modulation frequency ν_m is observed. By sweeping the modulation sideband through the Brillouin gain spectrum, we can resolve discrete acoustic modes of the fiber. The linewidth of the side-mode is measured to be less than 30 kHz and is limited by the resolution bandwidth of the spectrum analyzer and the linewidth of the microwave source as verified by comparing the side-mode and the microwave reference directly from microwave source (see inset).

We note that there are two dominant sidebands in a phase modulation (± 1 sideband). In this letter, we refer to the higher frequency tone as the blue sideband, and the lower frequency tone as the red sideband. By tuning the probe laser, we can scan the Brillouin spectrum with the blue or the red sidebands. As will be seen in Figs. 3 and 4 the blue-sideband and red-sideband spectra are mirror images of each other.

Fig. 3 shows the measured Brillouin gain for the DSF using the blue sideband. The pump power was set at 13 dBm. The optical phase modulator had a bandwidth of 8 GHz and was driven by a sinusoidal signal whose frequency was varied across the phonon sideband, while its amplitude was kept at 10 dBm corresponding to a peak-to-peak amplitude of 2 V. The probe power was 2 dBm for Fig. 3(a) and 9 dBm for Fig. 3(b). As shown in Fig. 3(a), we observe two side lobes in addition to the main peak in the Brillouin gain spectrum. The separation between the peak

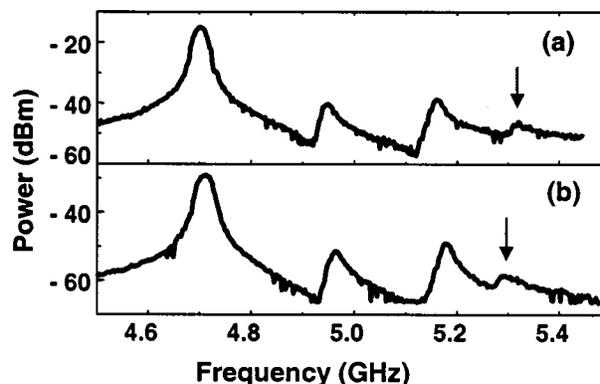


Fig. 4. Phonon spectra for two types of DSF. (a) With Corning SMF-DS. (b) With Corning SMF-LS.

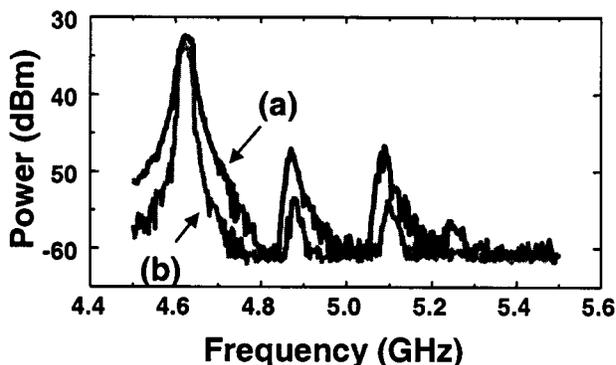


Fig. 5. Phonon spectrum of DSF with different microwave modulation power. (a) +20 dBm. (b) +10 dBm.

and the adjacent sidelobe is about 250 MHz, while the two lobes are apart by 210 MHz. These three peaks are very similar to the spectra reported in [9]. Comparing this with the calculated acoustic dispersion relation in [7], [9], we identify the peaks a, b and c with the L_{01} , $L_{02-L_{21}}$, and $L_{03-L_{22}}$ longitudinal modes, respectively. With a higher probe power in Fig. 3(b), we can clearly observe, for the first time, a fourth acoustic mode spaced by 150 MHz from $L_{03-L_{22}}$ mode. This mode, labeled as d, can be identified as $L_{04-L_{23}}$ mode with a power that is about 10 dB lower than the $L_{03-L_{22}}$ mode. The lower amplitude of this mode is due to the weaker overlap between the HE_{11} optical mode-field and the $L_{04-L_{23}}$ acoustic mode [12].

To further demonstrate the usefulness of the high-resolution microwave phonon spectroscopy technique, we characterized two different types of DSFs. We use Corning SMF-DS and SMF-LS fibers having Stokes frequencies (at 1550-nm pump wavelength) of 10.55 and 10.58 GHz, respectively. Fig. 4 shows the comparison of the longitudinal acoustic modes using the red sideband. The $L_{04-L_{23}}$ mode in SMF-LS is 100 MHz away from the $L_{03-L_{22}}$ mode. The three lower modes, namely the L_{01} , $L_{02-L_{21}}$, and $L_{03-L_{22}}$ longitudinal modes, are similar in these fibers. However, the L_{04} modes in these fibers are offset by about 50 MHz.

Fig. 5 shows the measured acoustic mode spectra with microwave input as a variable for the SMF-DS using the red sideband. As we increase microwave power from 10 to 20 dBm, we observe a 10-dB gain increment in the side-modes, while the L_{01} mode remains unchanged. In addition, the $L_{04-L_{23}}$ mode

peak rises above the measurement noise floor only at the higher microwave power. The dependence of Brillouin spectra on microwave power, leading to the observed phenomena, is under investigation. A possible cause is gain saturation in the Brillouin gain medium.

III. CONCLUSION

We report the first measurements of the acoustic mode spectrum of DSF. Using an improved technique to perform microwave phonon spectroscopy that provides high sensitivity and high spectral resolution, we report the first observation of the acoustic L_{04} mode in DSF. It is well known that the acoustic mode spectrum is sensitive to the fiber-doping profile and geometry. Therefore, high-resolution high-sensitivity Brillouin spectroscopy is a valuable tool that could be useful in development of new optical fibers.

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